

# Hyperfine Spectroscopy and Characterization of Muonium in ZnGeP<sub>2</sub>

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## The Experiment

• Transverse field (TF) hyperfine (HF) spectroscopy (Fig. 1) performed using *HiTime* high frequency spectrometer in the M15 beamline at TRIUMF, Vancouver Canada

• Longitudinal field (LF)  $T_1^{-1}$  depolarization measurements (Fig. 2) using EMU spectrometer at STFC ISIS Facility, UK

## Introduction

The purpose of this study is to extend the investigation of muonium defect centers as an experimentally accessible analog of isolated hydrogen defect states in semiconductors into the chalcopyrite structured II-IV-V<sub>2</sub> compounds. Also, we aspire to increase the understanding of probe dynamics associated with using positive muons to characterize the atomic-scale magnetism in Mn-doped II-IV-V<sub>2</sub> compounds.

The chalcopyrite structured II-IV-V<sub>2</sub> compounds are well developed and optimized as infrared active non-linear optical materials [1]. Several members of this class of compounds have shown room temperature ferromagnetism when doped with Mn, allowing these compounds to be possible candidates for future spintronics applications [2].

## The Sample

Shown (Right)

Nominally undoped single crystal of ZnGeP<sub>2</sub> with dimensions of 8x8x1 mm<sup>3</sup> and an [001] orientation

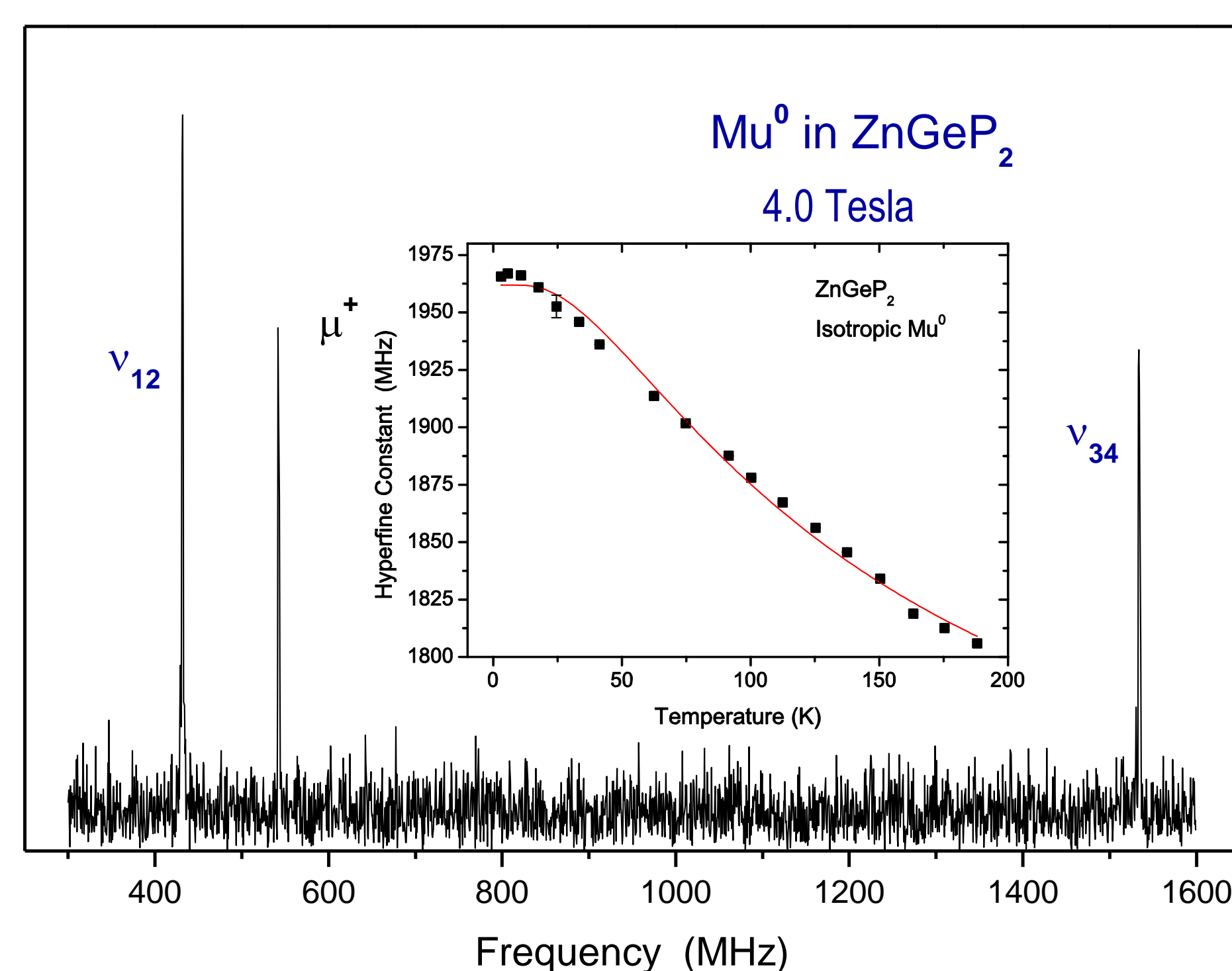
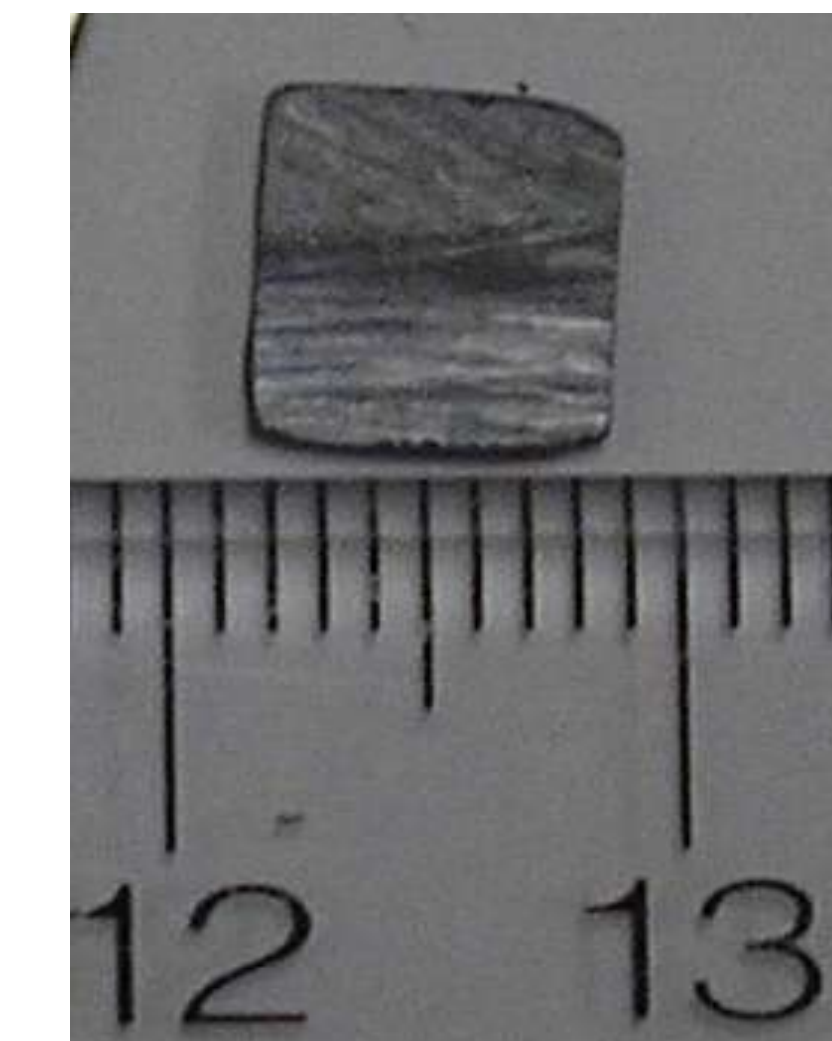


Fig. 1: Above is a TF MuSR FFT from the HF spectroscopy performed at  $T=5K$  and  $B=4.0T$ . The middle line, labeled  $\mu^+$ , at 542.14 MHz is the signal from the bare muon; where  $\nu_{12} = 432.1$  MHz (left) and  $\nu_{34} = 1534.8$  MHz (right) are the HF lines. Note that  $\nu_{12}$  precesses in a direction opposite of  $\mu^+$  and  $\nu_{34}$ . Inset shows variation of  $Mu^0$  hyperfine parameter as a function of temperature in  $B_{TF} = 4.0$  T. The solid line is the result of the current best model fit (Einstein model) following from equation 2 [3,4].

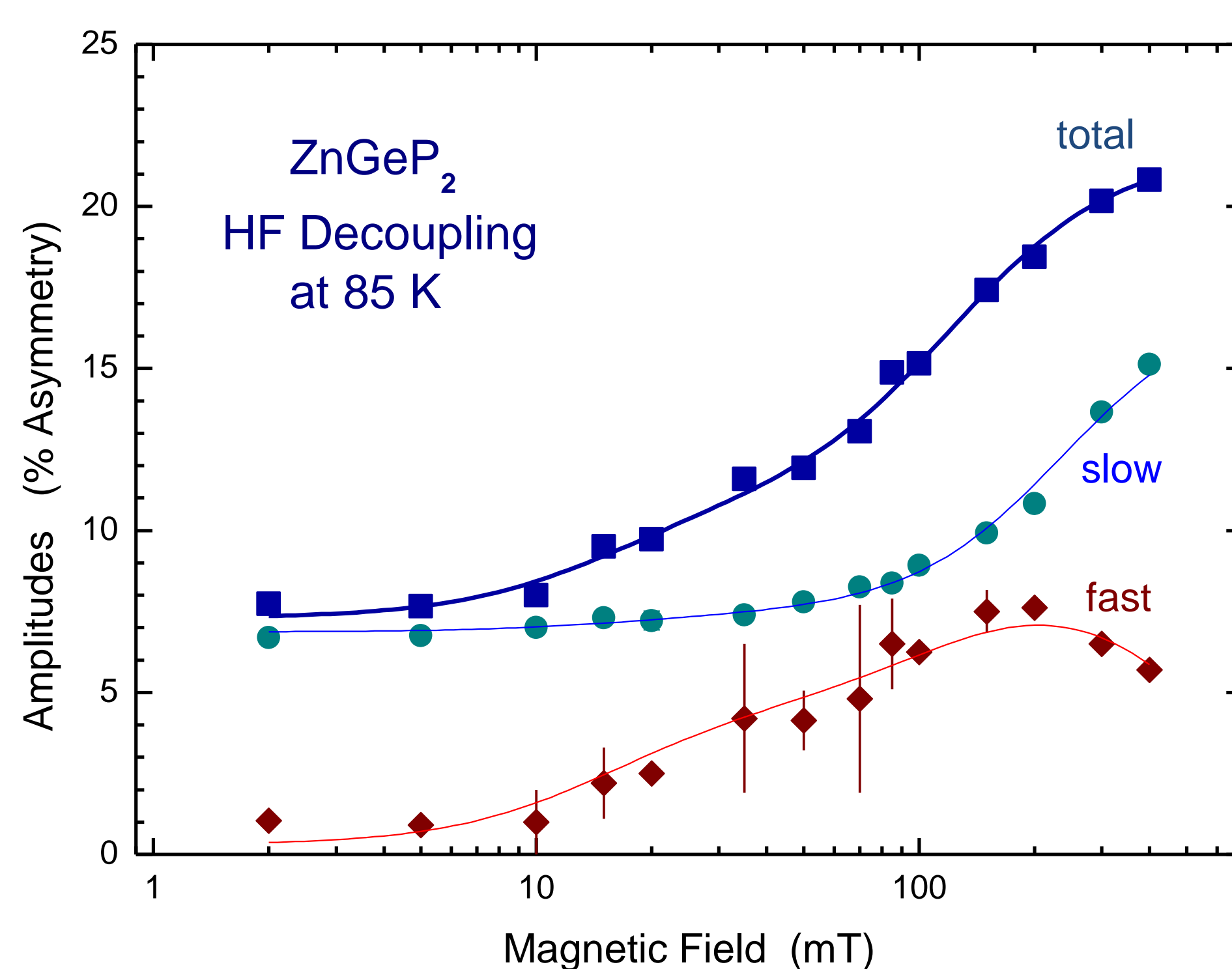


Fig. 2: Above is the HF decoupling curve from LF studies taken with fields from 0 to 400 mT. The flattening at both low and higher fields is indicative of an anisotropic hyperfine term. The isotropic piece having a value of roughly  $A_2 = 3185$  MHz and dipolar contribution of  $D = 374$  MHz.

## Results and Interpretation

- $\mu^+$  follows  $f = \gamma_\mu B$  where  $\gamma_\mu = 135.54$  MHz/T which is the diamagnetic frequency [4]
- $\nu_{12}$  and  $\nu_{34}$  show behavior characteristic of hyperfine splitting from a muonium atom,  $Mu^0$  with varying temperature and applied field [5]
- HF interaction for promptly formed  $Mu^0$  state (visible in TF) is isotropic since variations in sample orientation produced identical results.
- $T_1^{-1}$  depolarization data in LF, fig. 2, showed state with anisotropic hyperfine interaction [6]
- $T_1^{-1}$  depolarization rates as a function of field using single state model assuming dipolar fluctuating piece of the hyperfine parameter responsible for depolarization yields hop rates ( $1/\tau_c$ ), shown in Fig. 4, extracted from [7]

$$T_1^{-1} = \left(1 - x/\sqrt{1+x^2}\right) \frac{D^2 \tau_c}{1 + \omega_{12}^2 \tau_c^2} \quad (1)$$

where  $x = B_{LF}/B_2$  and  $B_2 = 0.1131$  T (decoupling field)

## Further Discussion

- Sites where  $Mu^0$  may reside:
  - T-site with Phosphorous near neighbors; where isotropic  $Mu^0$  formed
  - T-site associated with II-IV sublattice; where anisotropic  $Mu^0$  formed
- Variation in HF parameter, fig. 1 inset, as function of temperature fits local-mode vibrational model (Einstein model) [8,9] ie.
 
$$A(T) = A_0 + \frac{C_1}{\exp(h\nu/k_B T)} + \frac{C_2}{[\exp(k\nu/k_B T)]^2} \quad (2)$$

$A_0$  = isotropic zero-temperature HF parameter  
 $C_1$  &  $C_2$  = coupling constants (temperature independent)  
 $\nu$  = single vibrational frequency
- For this particular experiment,  $h\nu = 7.20$  meV, similar to  $Mu_{II}$  in structurally similar ZnSe [10]
- Hop rates and  $2\pi A_0$  having similar order of magnitude confirm the isotropic characteristic of the 1962 MHz HF term in T<sub>p</sub>-site and exclude the possibility of motional averaging

## Conclusions

- Upon muon implantation,  $Mu^0$  with isotropic hyperfine parameter of 1962 MHz, promptly formed in a phosphorous T-site
- $Mu^0$  rapidly hops, visiting other sites with different hyperfine parameters where there is an abrupt change in isotropic to anisotropic hyperfine parameter occurring at each subsequent hop
- While occupying the T-site associated with the II-IV sublattice,  $Mu^0$  has anisotropic hyperfine parameter (with axial symmetry) of  $A_2 = 3185$  MHz and  $D = 374$  MHz
- Since hop rates increase with decreasing temperature, motion below 80 K can be assigned to quantum tunneling; however, the slower decrease at lowest temperatures is commonly assigned to interactions with defects or impurities
- Motion between 100 K and 200 K is commonly assigned to thermally activated hopping
- Rapid increase in the relaxation rate (fig 3) above 200 K yields an ionization energy near 200 meV
- Further study necessary to develop a model for the relaxation due to motion involving two distinct  $Mu^0$  centers

## References:

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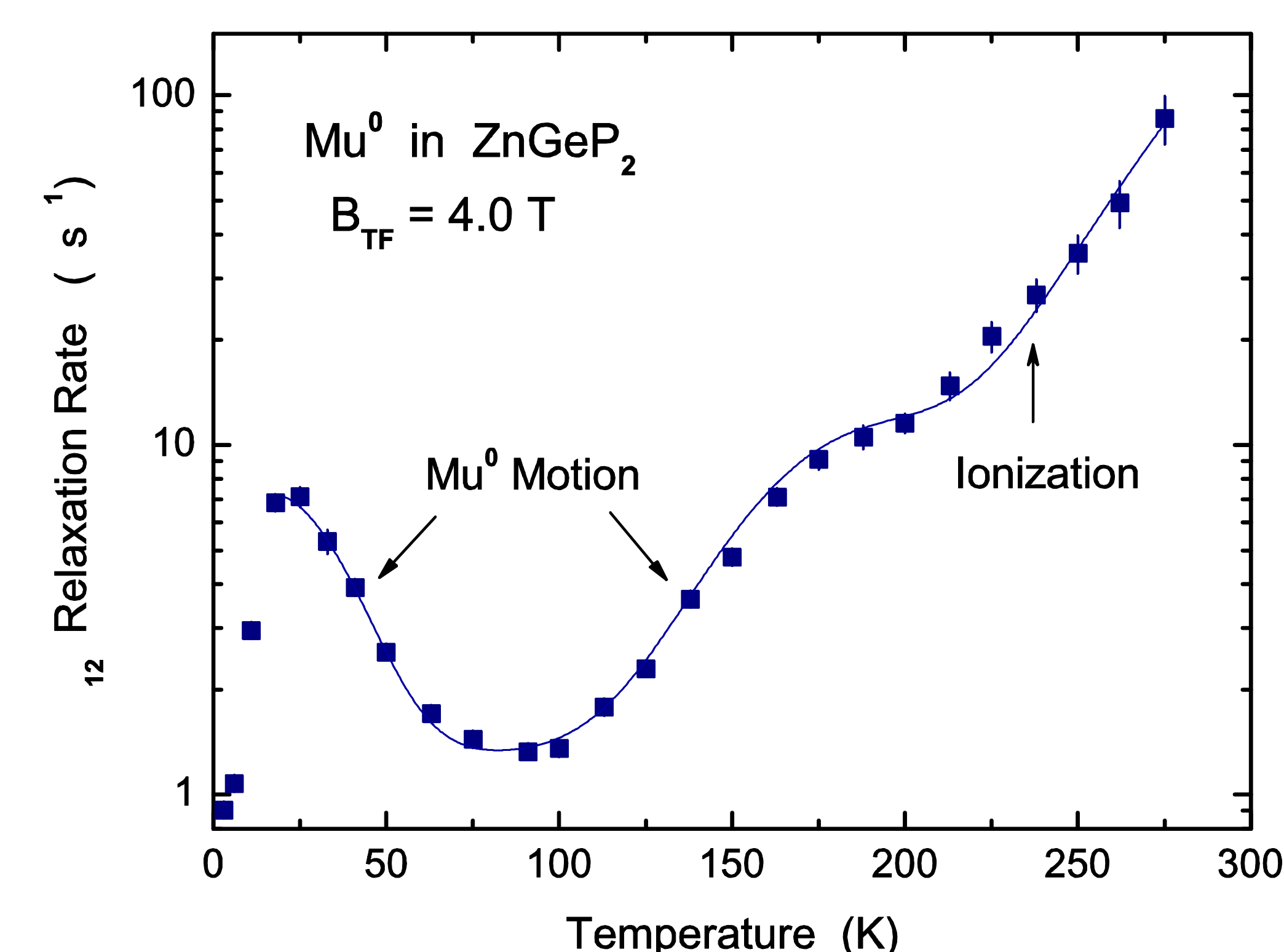


Fig. 3: Above is the relaxation rate of the  $\nu_{12}$  from TF measurements (LF shows similar results). The high relaxation rates at low temperatures ( $T < 50K$ ) have been assigned to quantum tunneling. Increasing relaxation rates from 100 to 200 K are assigned to thermally activated hops.

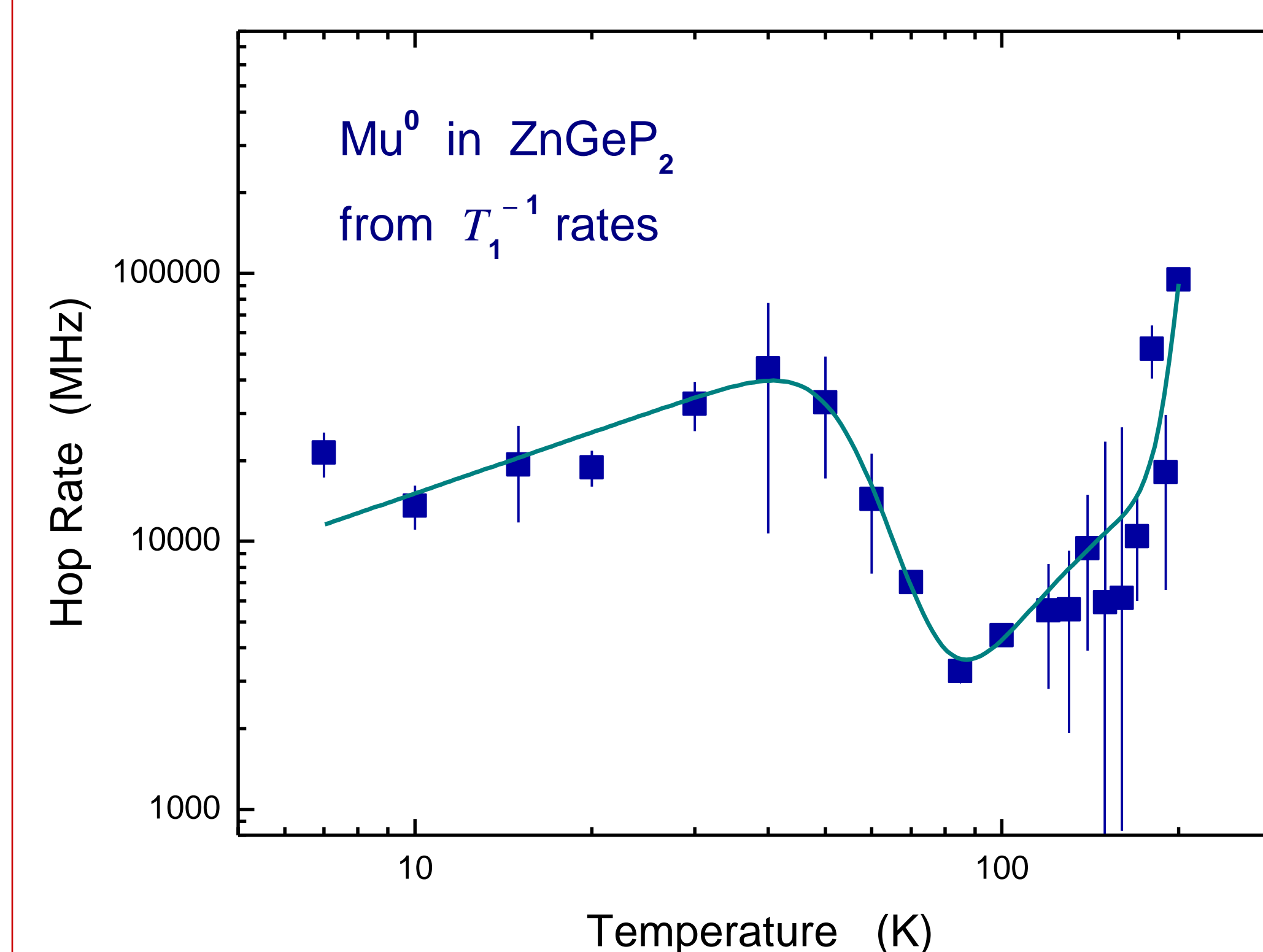


Fig. 4: Hop rates for anisotropic  $Mu^0$  in ZnGeP<sub>2</sub> as extracted from longitudinal depolarization data using equation 1.

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